## Lecture 11: Models for Pedestrian Flow

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### 11.1 Macroscopic Models for Unidirectional Flow

- Basic difference to (lane-based) vehicular traffic: full two-dimensionality
- Can also be applied to disordered traffic of other self-driven agents such as non-lanebased traffic flow, cycling, Marathon runs, inline-skating, and crosscountry-ski events
- In contrast to microscopic pedestrian models such as the Social-Force Model, macroscopic pedestrian models are only suited for essentially unidirectional pedestrian flows
- Because of the two dimensions, we need to redefine density and introduce a new macroscopic quantity: flow density



### 11.1.1 Elementary Macroscopic Variables of Twodimensional Flow

 y

- 2d density: $\rho(x, y, t)=\rho(\boldsymbol{x}, t)$ [pedestrians $/ \mathrm{m}^{2}$ or peds $/ \mathrm{m}^{2}$ ],
- Local velocity: $\boldsymbol{V}(\boldsymbol{x}, t)=\left(V, V_{y}\right)[\mathrm{m} / \mathrm{s}]$.
- Flow density: $\boldsymbol{J}(\boldsymbol{x}, t)=\left(J, J_{y}\right)=\rho \boldsymbol{V}$ [peds $/(\mathrm{ms})$ ]
- Effective 1d density [peds/m]:

$$
\rho_{1 \mathrm{~d}}(x, t)=\int_{y=-W / 2}^{W / 2} \rho(x, y) \mathrm{d} y \approx \rho W
$$

- Flow [peds/s]

$$
Q=\int_{y=-W / 2}^{W / 2} J(x, y) \mathrm{d} y \approx \rho V W \approx \rho_{1 \mathrm{~d}} V
$$

## Hydrodynamic relation and continuity equation

- hydrodynamic relations (vectorial and effective):

$$
\boldsymbol{J}=\rho \boldsymbol{V}, \quad J=\rho V, \quad Q=\rho_{1 \mathrm{~d}} V=\rho W V
$$

- Vectorial continuity equation:

$$
\frac{\partial \rho}{\partial t}+\frac{\partial J_{x}}{\partial x}+\frac{\partial J_{y}}{\partial y}=0
$$

- Effective continuity equation:

$$
\frac{\partial \rho_{1 \mathrm{~d}}}{\partial t}+\frac{\partial Q}{\partial x}=0
$$

important:

- In the effective continuity equation, lateral motion is introduced implictly: For example, in a stationary situation, $\frac{\partial}{\partial t}=0$, we have $Q=\rho W V=$ const.. If $W(x)$ narrows funnel-like, the 2 d -density $\rho$ increases by concentric lateral motion


### 11.1.2 Fundamental Diagramm




Parabolic FD a la Greenshields:

$$
\begin{aligned}
& V_{e}(\rho)=V_{0}\left(1-\frac{\rho}{\rho_{\max }}\right), \\
& J_{e}(\rho)=V_{0} \rho\left(1-\frac{\rho}{\rho_{\max }}\right)
\end{aligned}
$$

Weidmann FD:

$$
\begin{aligned}
v_{e}(\rho) & =v_{0}\left\{1-\exp \left[-\lambda\left(\frac{1}{\rho}-\frac{1}{\rho_{\max }}\right)\right]\right\} \\
J_{e}(\rho) & =\rho V_{e}(\rho)
\end{aligned}
$$

## Wave velocities



- Wave velocity $w(\rho)=J^{\prime}(\rho)$
- Notice: derivative with respect to 2 d density. Double density means $1 / \sqrt{2}$ times the average distance but more interacting people


### 11.1.3 Application I: Loveparade in Duisburg



Example for macroscopic event and evacuation planning:

- Assume unidirectional flow and a capacity density of $J_{\max }=1 \mathrm{ped} / \mathrm{s} / \mathrm{m}$ (a little bit lower than that of the Weidmann FD)
- Identify the bottleneck: The ramp to the event site at a width $W=30 \mathrm{~m}$ (the two tunnels have a summed cross section of $W=40 \mathrm{~m}$ )
- Calculate the bottleneck strength assuming no further obstacles:

$$
C=W J_{\max }=30 \mathrm{ped} / \mathrm{m}=108000 \mathrm{ped} / \mathrm{h}
$$

- Best case for three hours approach time (continuous unidirectional flow, no obstacles):

$$
n=324000 \text { persons }
$$

## Application II: Mixed traffic



## Application III: planning Marathon sports events



Planning of a starting scheme

- Define several starting groups $i$ with a maximum number $n_{i}$ of athletes, each, ordered to expected performance (best are first)
- Define a time delay $\tau$ between the starting shots of every group ("wave start"; the individual athlete's time starts when passing an electronic RFID gate)
- Identify the bottlenecks (often near the start) and estimate the flow profile from the expected speeds, their dispersion, and the start time delays
- Check if the maximum flow is below the bottleneck capacity; otherwise, change $n_{i}$ and $\tau$


## Application IV: Large-scale pedestrian streams in Mekka



## Application IV: Large-scale pedestrian streams in Mekka



### 11.2 Microscopic Model I: Social-Force Model

- In contrast to macroscopic models, microscopic pedestrian models are suited to model any pedestrian motion, whether directed or not:
- Any pedestrian can have its individual destination
- The models of this class are fully twodimensional; all pedestrians can move anywhere inside allowed regions
- The first and most prominent representative is the Social Force Model (SFM)
- More generally, the SFM is a model for self-driven particles or active particles, sometimes also called agents (no stirring or shaking involved)
- In analogy to Newtonian forces, the SFM pedestrian is driven by social forces
- In extended models, additional physical forces are modelled in case of a direct contact. However, this makes the equations hard to solve because it entails stiff differential equations (involving several time scales)


## Overall specification of the Social-Force Model

Pedestrians are active particles of mass 1, i.e., force equals acceleration:

$$
\dot{\boldsymbol{v}}_{i} \equiv \frac{\mathrm{~d} \boldsymbol{v}_{i}}{\mathrm{~d} t}=\boldsymbol{f}_{i}^{\mathrm{free}}+\sum_{j \neq i} \boldsymbol{f}_{i j}^{\mathrm{int}}+\sum_{k} \boldsymbol{f}_{i k}^{\mathrm{walls}}
$$

- The free force $f_{i}^{\text {free }}$ directs the pedestrians to their respective destinations. It can be modelled by a gradient of a free potential, e.g., indicating the shortest distance to the destination
- The pedestrian-pedestrian interaction forces $\boldsymbol{f}_{i j}^{\text {int }}$ acting from pedestrian $j$ onto $i$ are generally repulsive to avoid collisions and depends on the distance, directions, and the velocity vectors of both pedestrians
- The wall forces keep the pedestrians on the walkable area, i.e., away from the boundaries or from fixed compact obstacles. Both can be modelled as an obstacle floor field
- The floor fields and free potentials are static scalar fields $\Phi(x, y)$ defined on the walkable area. They can be pre-calculated for all boundaries, obstacles, and possible destinations


### 11.2.1 Free potential and force

- Free force corresponds to a vectorial free-traffic OVM

$$
\boldsymbol{f}_{i}^{\text {free }}=\frac{\boldsymbol{v}_{0 i}-\boldsymbol{v}_{i}}{\tau}
$$

- The desired velocity $\boldsymbol{v}_{0 i}=v_{0} \boldsymbol{\nabla} \Phi^{\text {free }}$ is given by the gradient $\boldsymbol{\nabla}=(\partial x, \partial y)$ of the free potential
- The free potential $\Phi^{\text {free }}(x, y)$ gives the distance to the destination on the shortest possible path
- Model parameters: speed adaptation time $\tau$ and desired speed (magnitude of $\boldsymbol{v}_{0 i}$ ) Give plausible values for $\tau$ and $v_{0}$


## Changing the pedestrian's preferences



- If the preference is to avoid "pancake" like crowding to leave doors etc, change the floor field to no longer reflect directly the shortest distance
- Here pedestrians prefer a queue rather than a "pancake"
- Rather than the gradient, just take $v_{0}$ times the unit vector of the gradient (Why was the gradient always a unit vector, in the last slide?)


### 11.2.2 Pedestrian-pedestrian interactions

Assumptions:

- The forces should be repulsive. ("Mexican-hat" potentials modelling attractive forces for couples, families, or friends will not be considered but can be added easily.)
- Generally, the forces should depend on the velocity vectors (speeds and directions) of both pedestrians. This also includes simple anticipation heuristics over anticipation time $\tau_{a}$
- Unlike physical forces, there is no momentum conservation (actio=reactio). Instead, forces from objects in viewing direction are stronger than that the nearly vanishing ones on the back. This anisotropy is the basis for fundamental diagrams
? Compare with car-following models
? What would a fundamental diagram look like for interaction forces satisfying actio=reactio?
$\Rightarrow$ general structure for a force from pedestrian/vehicle/compact obstacle $j$ onto pedestrian $i$ :

$$
\boldsymbol{f}_{i j}=\boldsymbol{f}_{i j}\left(\boldsymbol{x}_{i}, \boldsymbol{x}_{j}, \boldsymbol{v}_{i}, \boldsymbol{v}_{j} ; \tau_{a}\right)
$$

## Further simplifying assumptions for the interactions

The general structure implies two functions $f_{x}($.$) and f_{y}($.$) with eight dynamic arguments$ (the components of the two position and velocity vectors), each.
Rather complicated $\Rightarrow$ Simplify further

- Translational invariance: Dependence on the distance vector $d_{i j}=x_{i}-x_{j}$, only, and not separately on $\boldsymbol{x}_{i}$ and $\boldsymbol{x}_{j}$ in difference vectors, the subject always comes first!
- Without the viewing angle dependency, the forces can be written as a gradient of an interaction potential, $\boldsymbol{f}_{\text {pot }}=\nabla_{\boldsymbol{d}} \Phi^{\text {int }}\left(\boldsymbol{d}, \boldsymbol{v}_{i}-\boldsymbol{v}_{j}\right)$
- The potential/potential force is either Galilei invariant, i.e., depends on the relative velocity vector $v_{i}-v_{j}$, only (Elliptical specification II), or the velocity of the interacting pedestrian/object $j$ is ignored (Elliptical specification I), or there is neither velocity dependence nor anticipation (Circular specification). (in any case, the free part and the obstacles break this invariance on the system level)
- The viewing angle dependency is just a multiplicative prefactor $w(\cos \phi)$ of the cosine of the viewing angle, $\cos \phi=-\boldsymbol{d} \cdot \boldsymbol{v}_{i} /\left(|\boldsymbol{d}|\left|\boldsymbol{v}_{i}\right|\right)$
$\Rightarrow$ Force expression reduced to a scalar potential of only four dynamical variables:

$$
\boldsymbol{f}_{i j}=w(\cos \phi) \boldsymbol{f}_{i j}^{\mathrm{pot}}\left(\boldsymbol{d}, \boldsymbol{v}_{i}-\boldsymbol{v}_{j}\right)=-w(\cos \phi) \nabla_{\boldsymbol{d}} \Phi^{\mathrm{int}}\left(\boldsymbol{d}, \boldsymbol{v}_{i}-\boldsymbol{v}_{j}\right)
$$



## Anticipation

Pedestrian 2
(forces not considered)


## Constructing the interaction potential for the elliptical specifications



- Determine the present collision point $F_{1}=\boldsymbol{x}_{j}$ for $i$ and the present location $F_{2}=\boldsymbol{x}_{j}-\Delta \boldsymbol{d}$ for an anticipated collision after $\tau_{a}$
- Define ellipses by the focal points F1 and F2 and the present distance vector $\boldsymbol{d}$
- Equipotential lines have a constant semi-minor axis $b$
- Hammer two nails at F1 and F2 and attach the ends of a string of length $L>\overline{F_{1} F_{2}}$ to the nails. Tighten the string with a pencil and draw. You will draw an ellipse which therefore satisfies


## Constructing the interaction potential for the elliptical specifications



This also applies to point $C$ defining the minor semi-axis $b$ of the ellipse:

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$$
\begin{gathered}
\overline{F_{1} P}+\overline{P F_{2}}=L \\
|\boldsymbol{d}|+|\boldsymbol{d}+\Delta \boldsymbol{d}|=L
\end{gathered}
$$

## Constructing the interaction potential for the elliptical specifications



This also applies to point $C$ defining the minor semi-axis $b$ of the ellipse:

$$
\begin{aligned}
|\boldsymbol{d}|+|\boldsymbol{d}+\Delta \boldsymbol{d}|=L & =2 \sqrt{\left(\frac{\mid \Delta \boldsymbol{d}}{2}\right)^{2}+b^{2}} \\
b(\boldsymbol{d}) & =\frac{1}{2} \sqrt{(|\boldsymbol{d}|+|\boldsymbol{d}+\Delta \boldsymbol{d}|)^{2}-|\Delta \boldsymbol{d}|^{2}}
\end{aligned}
$$

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$$
\begin{gathered}
\overline{F_{1} P}+\overline{P F_{2}}=L \\
|\boldsymbol{d}|+|\boldsymbol{d}+\Delta \boldsymbol{d}|=L
\end{gathered}
$$

## The SFM interaction potential

- The focal line $\overline{F_{1} F_{2}}$ defining generated ellipses with $b=0$ means collision with pedestrian/object $j$ within the anticipation horizon $\left[t, t+\tau_{a}\right]$
- $\Rightarrow$ potential should be highest for $b(x, y)=0$ and decrease with $b$ :

$$
\Phi^{\mathrm{int}}(\boldsymbol{d})=A B \exp \left(\frac{-b(\boldsymbol{d})}{B}\right), \quad b(\boldsymbol{d})=\frac{1}{2} \sqrt{(|\boldsymbol{d}|+|\boldsymbol{d}+\Delta \boldsymbol{d}|)^{2}-|\Delta \boldsymbol{d}|^{2}}
$$

- Potential depends on the present distance vector $\boldsymbol{d}=\boldsymbol{x}_{\boldsymbol{i}}-\boldsymbol{x}_{\boldsymbol{j}}$ and on the anticipated distance vector change
- Elliptical specification II: $\Delta \boldsymbol{d}=\tau_{a}\left(\boldsymbol{v}_{\boldsymbol{i}}-\boldsymbol{v}_{j}\right)$
- Elliptical specification I: $\Delta \boldsymbol{d}=\tau_{a} \boldsymbol{v}_{\boldsymbol{i}}$
- Circular specification: $\Delta \boldsymbol{d}=\mathbf{0}$
- parameters:
- interaction strength $\boldsymbol{A}$, typically values around $A=2 \mathrm{~m} / \mathrm{s}^{2}$
- range $B$, typically values around $B=1 \mathrm{~m}$
- For the circular definition without anticipation, we have $\Phi^{\text {int }}(\boldsymbol{d})=A B e^{-|d| / B}$


## Questions

? Potentials if both pdestrians have the same velocity

## Proposal of an improved and simpler potential: Anticipated circular specification

The elliptical specification II shows the most plausible behaviour and performs best in calibration/validation. Still it has some imperfections:

- Gradient, i.e., the derived social forces, diverge towards the focal points of the ellipse
- Ad-hoc nature. Why ellipses and the complicated construction?
- Answer: Use the circular potential but centered at a position where the two pedestrians come closest within the anticipation horizon


## Anticipated circular specification

1. Calculate the time interval $\tau^{\prime}$ of closest encounter within the anticipation time horizon:

- anticipated distance vector $\boldsymbol{d}_{a}\left(\tau^{\prime}\right)=\boldsymbol{d}+\left(\boldsymbol{v}_{i}-\boldsymbol{v}_{j}\right) \tau^{\prime}$
- Minimum distance reached ( $\boldsymbol{d}_{a}$ perpendicular to $\boldsymbol{v}_{i}-\boldsymbol{v}_{j}$ ) at time

$$
\tau_{\min }=-\frac{\boldsymbol{d} \cdot\left(\boldsymbol{v}-\boldsymbol{v}_{j}\right)}{\left(\boldsymbol{v}-\boldsymbol{v}_{j}\right)^{2}}
$$

- anticipated time

$$
\tau^{\prime}=\max \left(0, \min \left(\tau_{a}, \tau_{\min }\right)\right)
$$

2. Construct circular potential around the distance vector at $t^{\prime}$ :

$$
\Phi^{\mathrm{int}, \bmod }(\boldsymbol{d})=A B \exp \left(-\frac{\left|\boldsymbol{d}^{\prime}\right|}{B}\right), \quad \boldsymbol{d}^{\prime}=\boldsymbol{d}+\left(\boldsymbol{v}_{i}-\boldsymbol{v}_{j}\right) \tau^{\prime}
$$

- If the time of shortest encounter lies in the past, we have just the normal circular potential
- Since the potential center generally is different for each point ( $x, y$ ), we, in effect, have an elliptical-like potential but without divergent gradients


## Comparison of potentials



Detail of original (left) and modified (right) interaction potentials for $\boldsymbol{v}_{1}=(1,0)$ and $\boldsymbol{v}_{2}=\mathbf{0}$ near the focal point $F_{1}=(0,0)$ of the original SFM.

## SFM potential and forces on a pedestrian walking in NE direction ...



- caused by a standing or moving pedestrian/compact object (circular specification)
- In the circular specification, the own velocity only influences the directional dependence (not shown here). The velocity of the interacting pedestrian is ignored


## SFM potential and forces on a pedestrian walking in NE direction ...



- caused by a standing pedestrian/compact object (specification II)
- or by standing/moving objects (specification I)
- Note the different potentials for moving targets in spec I and II


## SFM potential and forces on a pedestrian walking in NE direction ...



## SFM potential and forces on a pedestrian walking in NE direction ...



- caused by a moving pedestrian (improved/modified specification II)
- Although circular for a given point $(x, y)$, it becomes elongated as a function of $(x, y)$ because the center (closest anticipated distance) changes with $(x, y)$


## Deriving the force field from the potential

- The potential $\Phi^{\text {int }}$ is de-facto a "hill" with a ridge along the positions where the subject pedestrian $i$ expects a collision within the anticipation horizon.
- The potential force $-\nabla \Phi^{\text {int }}$ is directed "downhill", i.e., away from potential collision points
- Use the fact, that, for arbitrary constant vectors $\boldsymbol{a}$, we have

$$
\nabla_{d}|d+a|=\frac{d+a}{|d+a|}=e_{d+a}
$$

## Circular potential

$$
\begin{aligned}
f^{\mathrm{pot}}=-\nabla_{d} \Phi^{\mathrm{int}}(\boldsymbol{d}) & =-\nabla_{d}\left(A B e^{-\frac{|d|}{B}}\right) \\
& =A e^{-\frac{d}{B}} \nabla_{d}|\boldsymbol{d}|
\end{aligned}
$$

$$
f^{\text {pot }}=A e^{-\frac{d}{B}} e_{d} \quad \text { circular potential }
$$

## Deriving the force field for the elliptical specifications I and II

The two speciifcations differ only in the constant $\Delta \boldsymbol{d}=\boldsymbol{v}_{i} \tau_{a}(\mathrm{I})$ and $\Delta \boldsymbol{d}=\left(\boldsymbol{v}_{i}-\boldsymbol{v}_{j}\right) \tau_{a}$, respectively

$$
\Phi^{\mathrm{int}}(\boldsymbol{d})=A B e^{-\frac{b(\boldsymbol{d})}{B}}, \quad b(\boldsymbol{d})=\frac{1}{2} \sqrt{(|\boldsymbol{d}|+|\boldsymbol{d}+\Delta \boldsymbol{d}|)^{2}-|\Delta \boldsymbol{d}|^{2}}
$$

$$
\begin{aligned}
f^{\text {pot }}=-\nabla_{d} \Phi^{\mathrm{int}}(\boldsymbol{d}) & =-\nabla_{d}\left(A B e^{-\frac{b(\boldsymbol{d})}{B}}\right) \\
& = \\
& A e^{-\frac{d}{B}} \nabla_{d}(b(\boldsymbol{d})) \\
& =A e^{-\frac{d}{B}} \frac{|\boldsymbol{d}|+|\boldsymbol{d}+\Delta \boldsymbol{d}|}{4 b}\left(\boldsymbol{e}_{\boldsymbol{d}}+\boldsymbol{e}_{\boldsymbol{d}+\Delta \boldsymbol{d}}\right) \\
\text { thread and nails } & A e^{-\frac{d}{B}} \frac{1}{4} \sqrt{4+\left(\frac{\Delta \boldsymbol{d}}{b}\right)^{2}}\left(\boldsymbol{e}_{\boldsymbol{d}}+\boldsymbol{e}_{\boldsymbol{d}+\Delta \boldsymbol{d}}\right),
\end{aligned}
$$

$$
\boldsymbol{f}^{\text {pot }}=A e^{-\frac{b(\boldsymbol{d})}{B}} \sqrt{1+\left(\frac{\Delta \boldsymbol{d}}{2 b(\boldsymbol{d})}\right)^{2}}\left(\frac{\boldsymbol{e}_{\boldsymbol{d}}+\boldsymbol{e}_{\boldsymbol{d}+\Delta \boldsymbol{d}}}{2}\right) \quad \text { Elliptical specifications I and II }
$$

## Deriving the force field for the anticipated circular specification

This is like the potential for the circular specification, but at the anticipated position $\boldsymbol{d}+\Delta \boldsymbol{d}^{\prime}$ (shifted from $\boldsymbol{d}$ by a constant vector):

$$
\boldsymbol{f}^{\text {pot }}=A \exp \left(-\frac{\left|\boldsymbol{d}+\Delta \boldsymbol{d}^{\prime}\right|}{B}\right) \boldsymbol{e}_{d+\Delta \boldsymbol{d}^{\prime}} \quad \text { Anticipated circular specification }
$$

with

$$
\begin{aligned}
\Delta \boldsymbol{d}^{\prime} & =\left(\boldsymbol{v}_{i}-\boldsymbol{v}_{j}\right) \tau^{\prime} \\
\tau^{\prime} & =\max \left(0, \min \left(\tau_{a}, \tau_{\min }\right)\right) \\
\tau_{\min } & =-\frac{\boldsymbol{d} \cdot\left(\boldsymbol{v}_{i}-\boldsymbol{v}_{j}\right)}{\left(\boldsymbol{v}_{i}-\boldsymbol{v}_{j}\right)^{2}}
\end{aligned}
$$

- If $\boldsymbol{v}_{i}=\boldsymbol{v}_{j}$, both the anticipated circular potential and the elliptical specification II revert to the circular potential but not the elliptical specification I. (why?)


## Directionality

The potential force ignores that forces from pdestrians/objects ahead are stronger than that in the back and that also relative speed should matter. Multiplicative approach:

$$
\boldsymbol{f}^{\mathrm{int}}=w \boldsymbol{f}^{\mathrm{pot}}
$$

(1) Classical SFM dependency on the viewing angle:

$$
w(\cos \phi)=\lambda+(1-\lambda)\left(\frac{1+\cos \phi}{2}\right), \quad \cos \phi=-\boldsymbol{e}_{v_{i}} \cdot \boldsymbol{e}_{d}
$$




## Directionality (FVDM approach)

(2) Dependence on the approaching rate (generalisation of the FVDM relative speed sensitivity):

$$
\begin{equation*}
w(\dot{d})=\max (0,1-\gamma \dot{d}), \quad \dot{d}=\boldsymbol{e}_{d} \cdot\left(\boldsymbol{v}_{i}-\boldsymbol{v}_{j}\right) \tag{1}
\end{equation*}
$$



From a symmetric to a directed asymmetric interaction: Pedestrian at $(x, y)$ walking to a target pedestrian moving to the right




## From the asymmetric interaction to the full force: adding the free force




## Viewing angle vs. approaching rate directional weighting (example anticipated circular potential)




### 11.2.3 Potential and force by obstacles

- Basically, obstacles and boundaries of walkable areas are like standing pedestrians: One should not collid with/transgress them
- Small compact obstacles (poles, signposts, trees, pillars) can, in fact, be handled like standing pedestrians
- For extended obstacles/boundaries this would be inefficient (a cordon of many standing virtual pedestrians) and also biased (the additive effect of the forces exaggerates the effect)


## ? What to do ?

! Use the fact that obstacles, boundaries etc are really immobile so as to precalculate a global floor potential from all obstacles and boundaries
! Add anticipation in the same way as for defining the anticipated circular potential: Calculate the anticipation point and take the gradient at that point
! Multiply a factor for the viewing angle and/or approaching rate as when interacting with other pedestrians

## Potential and force by obstacles: precalculate the obstacle floor field



- Identify all obstacles and boundary objects
- Define a grid on the walkable area. For each gridpoint (at $\boldsymbol{x}$ ) and each obstacle object $k$, determine the distance vector $\boldsymbol{d}_{k}(\boldsymbol{x})$ to the neares point of this object
- Ignore obstacles that are shielded or too far away
- The global floor potential is given by

$$
\Phi^{\mathrm{obs}}(\boldsymbol{x})=A B \sum_{k} e^{-\frac{d_{k}(\boldsymbol{x})}{B}}
$$

## Obstacle social forces on pedestrians

The procedure for a pedestrian at $\boldsymbol{x}_{i}$ and velocity $\boldsymbol{v}_{i}$ is the same as for the anticipated circular interaction potential:

- Calculate the anticipated position $\boldsymbol{x}_{i}^{\prime}$ along the line $\left\{\boldsymbol{x}_{i}+t^{\prime} \boldsymbol{v}_{i}: t^{\prime} \in\left[0 . \tau_{a}\right]\right\}$
with the shortest distance $\boldsymbol{d}^{\prime}$ to

any obstacle (or with the highest precalculated gradient)
- The obstacle social force, i.e., acceleration, is given by

$$
\boldsymbol{f}_{i}^{\mathrm{obs}}=-w\left(\cos \phi_{i}^{\mathrm{obs}}\right) \nabla_{x} \Phi^{\mathrm{obs}}\left(\boldsymbol{x}_{i}^{\prime}\right)
$$

- A bilinear interpolation on the precalculated grid is enough for calculating the gradient
- If only a single obstacle is relevant, we have the anticipated circular specification:

$$
\boldsymbol{f}^{\mathrm{obs}}=w\left(\cos \phi_{i}^{\mathrm{obs}}\right) A e^{-\frac{d^{\prime}}{B}} \boldsymbol{e}_{d^{\prime}}
$$

### 11.2.4 Model Parameters and Fundamental Diagram

Overview of the model parameters of the SFM

| Parameter | normal <br> walking | Marathon <br> runners | comment |
| :--- | :--- | :--- | :--- |
| Desired speed $v_{0}$ <br> Speed adaptation time $\tau$ | $1.2 \mathrm{~m} / \mathrm{s}$ <br> Interaction range $B$ | $3 \mathrm{~m} / \mathrm{s}$ <br> 1.5 s | free traffic <br> free traffic |
| Interaction strength $A$ | $2 \mathrm{~m} / \mathrm{s}^{2}$ | $3 \mathrm{~m} / \mathrm{s}^{2}$ | of the order of the maximum <br> acceleration <br> decays by a factor $1 / e$ per <br> distance increment $B$ <br> anticipation for collisions <br> assuming constant velocities |
| directionality $\lambda$ | 1 s | 2 s | 0.03 |
| relative speed sensitivity $\gamma$ | 1.5 | 1.0 | isotropic actio=reactio: $\lambda=1$ <br> alternative formulation of the <br> directionality (as in the FVDM) |

## Fundamental diagram

- Because of the different possible geometric configurations, a fully 2d fundamental diagram (FD) is not unique
- A simpler approach is to define a single-file fundamental diagram
- Because the number of interacting persons in single files increases linearly with distance rather than quadratically, a single-file FD as a function of the 1 d density $\rho_{1 d}$ also approximates a 2 d FD as a function of the 2 d density $\rho$
As usual in FDs, we have identical pedestrians with identical (center-center) distances $\Delta x=1 / \rho_{1 d}$ and identical speeds



## Derivation for full interactions without shielding

$$
\begin{aligned}
& \frac{\mathrm{d} v_{i}}{\mathrm{~d} t}=\frac{v_{0}-v_{i}}{\tau}+\sum_{l=1}^{\infty} f_{i l}+\sum_{m=1}^{\infty} f_{i m} \\
&=\quad \frac{v_{0}-v_{i}}{\tau}-1 \sum_{l=1}^{\infty} A e^{-l \Delta x / B}+\lambda \sum_{m=1}^{\infty} A e^{-m \Delta x / B} \\
&=\quad \frac{v_{0}-v_{i}}{\tau}-A(1-\lambda) \sum_{l=1}^{\infty} e^{-l \Delta x / B} \\
&=\frac{v_{0}-v_{i}}{\tau}-A(1-\lambda)\left(\sum_{l=0}^{\infty} e^{-l \Delta x / B}-1\right) \\
& \text { geometric series } \frac{v_{0}-v_{i}}{\tau}-A(1-\lambda)\left(\frac{1}{\left.1-e^{-\Delta x / B}-1\right)!} \stackrel{!}{=} 0\right. \\
& v_{i}(\Delta x) \rightarrow V(\Delta x)=v_{0}-\tau A(1-\lambda)\left(\frac{e^{-\Delta x / B}}{1-e^{-\Delta x / B}}\right)
\end{aligned}
$$

## SFM fundamental diagrams for a single file without shielding


$Q\left(\rho_{1 \mathrm{~d}}\right)=\rho_{1 \mathrm{~d}} V\left(1 / \rho_{1 \mathrm{~d}}\right), \quad V(\Delta x)=v_{0}-\tau A(1-\lambda)\left(\frac{e^{-\Delta x / B}}{1-e^{-\Delta x / B}}\right)$

## SFM FD for a single file with NN interactins (one front and back pedestrian)



The value of $A$ such that $V_{\mathrm{NN}}(0.5 \mathrm{~m})=0$
$V_{\mathrm{NN}}(\Delta x)$ corresponds to third line of the derivation with only $l=1$ :

$$
Q\left(\rho_{1 \mathrm{~d}}\right)=\rho_{1 \mathrm{~d}} V_{\mathrm{NN}}\left(1 / \rho_{1 \mathrm{~d}}\right), \quad V_{\mathrm{NN}}(\Delta x)=v_{0}-\tau A(1-\lambda) e^{-\Delta x / B}
$$

## 11.3: An alternative approach: PLEdestrian

